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# Electrical and photoelectric properties of crystal InGaTe<sub>2</sub>

### ABSTRACT

Electroconduction, photoconduction and the curves of relaxation of notequilibrium photoconduction in the crystals  $InGaTe_2$  under the influence of laser and collinear light were carried out experimentally. It is shown that the optical absorption in  $InGaTe_2$  is carried out by indirect and direct optical transitions, the band gap is accordingly 1,02eV and 1,42eV. Acceptor centers are with the depth of 0.203 eV and 0.801 eV. At high levels optical excitation in  $InGaTe_2$  observed two-photon absorption. It was found that heat non-equilibrium photoconduction detected in the crystals  $InGaTe_2$  under the first and second harmonic of radiation of Neodymium laser caused impurity, bipolar and two-photon photoconduction.

Keywords: compound İnGaTe<sub>2</sub>, indirect and direct optical transitions, non-linear optic properties,
 photoelectric properties

### 1. INTRODUCTION

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Complex semiconductor compounds of the class  $A^{III}B^{III}C_2^{VI}$  are interesting as a material for 20 recieves of radiation. Radiation resistance and effect of radiation on electrical conductivity is intensively 21 22 studied in compounds as  $TlGa(In)Te_{2}(Se,S)$  [1-6]. The crystal InGaTe, being the representatives of the 23 mentioned class of compounds, have not been studied enough compared with other compounds. The power zonal chart was calculated theoretically by the pseudopotential method [7] and dispersion curves of 24 25 high frequency dielectric permittivity and temperature dependences of differential termo-emf [8,9] were studied. The obtained date are not sufficient for making general idea on physical properties of monocrystal 26 27  $InGaTe_{2}$  and as well as for specifying values of the fundamental parameters.

28 In this work presented experimental results of the investigation of electroconductivity and 29 photoconductivity of  $InGaTe_2$  width range of temperature of different values of optic excitation.

As is known, at higher degrees of optic excitation a sufficiently high concentration nonequilibrium carrier of current is achieved and this involves emergency of certain qualitatively new features of semiconductors as non-linear optic and photoelectric properties, saturation effects at absoption of light, etc. The use of laser radiation with tunable wavelength, and with different excitation intensities can simultaneously detect impurity, intrinsic, as well as nonlinear photoconductivity, if any, are in the investigated crystals. Experimental investigation of generation recombination processes in  $InGaTe_2$  laser radiation is of great interest because of their application in optoelectronics.

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## 2. METHODS OF CALCULATION

The single crystals  $InGaTe_{2}$  were grown by the Bridgman method have tetragonal crystal structure with spatial symmetry  $D_{4h}^{18}(I4/mcm)$ , with lattice parameters  $a = 8,463 \stackrel{0}{A}$  and  $c = 6,981 \stackrel{0}{A}$ . The Bridgman

42 spatial symmetry  $D_{4h}^{18}(I4/mcm)$ , with lattice parameters a = 8,463A and c = 6,981A. The Bridgman 43 method is a popular way of producing semiconductor crystals. The samples for measurements were 44 cut out from an ingot in the form of a rectangle with the sizes of  $6 \times 3 \times 1$  mm<sup>3</sup> and with such an orientation 45 that the lighting happens in the direction parallel to the crystallography axes "c". The monocrytal 46 InGaTe 2 possesed p-type conductivity and depending on the mode of a method of cultivation of a

crystal, specific resistand varied at the range  $(1 \cdot 10^3 \div 1 \cdot 10^6)$  Om·sm. A silver paste is used as ohmic 47 contact . Photoconductivity of the crystal InGaTe , was investigated by collinear source of light at 48 stationary mode, by the method of modultion of light intensity at frequency at 47kHz. At higher level of 49 50 optic excitation, the pulse laser Nd: YAG with built in generators of the second and third harmonic intended for generation of radiation with wavelength 1064, 532, 355nm with reconstructed wavelength at 51 the range from 410-710nm was used as radiation source the impulse duration  $\Delta t = 1.10^{-8}$  sec; maximum 52 power ~12 MVt/sm<sup>2</sup>. Intensity of laser radiation was measured (changed) by means of calibrated neutral 53 filters. In the work nonstationary digital system including memorable oscillograph and computer system 54 55 was used. The experimental technique is similar to one described in the paper [10].

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#### 3. RESULTS AND DISCUSSION

Dependence of electrical conductivity  $\sigma$  on temperature T is represented in Fig.1. As is seen from the figure, in the dependence  $\lg \sigma \sim 10^3 / T$  one can isolate two rectilinear areas belonging to generation of charge carrier with electronic transitions  $E_1 = 0,203$  3B and  $E_2 = 0,801$  eV. At wide interval of temperature 100-500 K the samples have p-type conductivity, we can note that temperature dependence of  $\sigma$  is stimulated by thermal ionization of acceptor centers and activation of electrones from the filled acceptor level. It should be noted that the temperature dependence of the electrical conductivity in different samples are identical.

66 Photoconductivity of the crystal  $InGaTe_2$  were investigated at interval of temperature 132–373K 67 Fig. 2. Photocurrent range cover a very wide range of energy from 1,2eV to 2,6eV. Photosensitivity of the 68 crystals steadily increases according to temperature increase. As is seen from the figure the long-wave 69 edge of ranges of photocurrent don't subject to Moss criterion. The character optic transition in InGaTe<sub>2</sub> is

seen from dependence  $I_{\phi}(h\nu)$ . In the range  $I_{\phi}^{\frac{1}{2}} \sim h\nu$  experimental points in the area of long-wave

edge lie on the same line that indicates the availability of indirect transition with energy 1,02eV (fig.2a).
Determination of the width of the band structure is of semiconductors by this method is quite justified . As
absorption coefficient for indirect transitions is much less than for direct transitions, its definition from the
absorption range is not really a simple task. Many authors use for that a more sensitive method as the
photoconductivity method.

76 It is known that the absorption coefficient for indirect transition for equals

$$\alpha \sim (h\nu - E_g^{\rm in} \pm E_p)^2 \tag{1},$$

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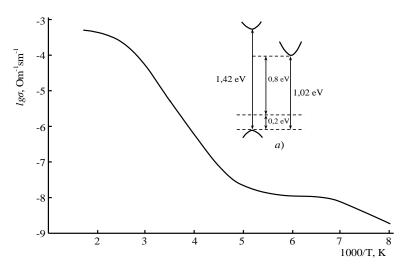


Fig. 1. The dependence of the electrical conductivity of crystal  $InGaTe_2$  on the temperature. a) - the energy diagram  $InGaTe_2$ 

where hv- is the energy of quantum,  $E_g^{\rm in}$  - is the width of the band structure,  $E_p$  -is the energy of the phonon that participating in indirect optic transitions. As the photoconductivity size  $I_F^{1/2} \sim \alpha^{1/2}$ , then removing the dependence of  $I_F$  on energy  $h\nu$  in the weak abroption zone (at energies less than the direct zone) one can determine the width of the band structure of indirect transitions.

83 Maximum photosensitivity of the crystal  $InGaTe_2$  is achieved at 1,42eV. High photoconductivity 84 in the fundamental absorption area testified on availability of direct optical transitions in the crystal. 85 Therefore, we can suppose that, in addiction to indirect transition in the monocrystal  $InGaTe_2$ , the direct 86 optical transition with the energy 1,42eV stipulated generation of nonequilibrium charge carriers.

87 Temerature dependence of photocurrent in  $InGaTe_2$  is represented in fig.3. As is seen from the Fig. 3,

temperature clearing of photocurrent happens at higher than 373K. The range of photocurrent remined in

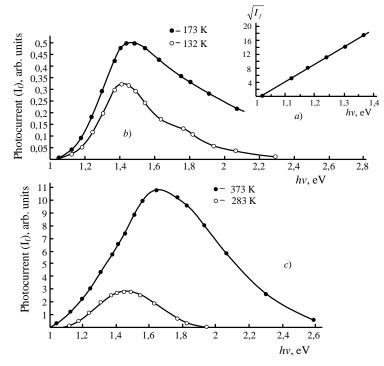


Fig. 2. Photoconductivity spectra of crystal  $InGaTe_2$  with different temperatures.

- the temperature clearing area has a sharp edge in the area of long-wave absorption and maximum of the
- 90 range is at 1,42eV. It is clear that in the area of temperature quenching of photocurrent main contribution
- 91 to the photoconductivity make nonequilibrium charge carriers generated as a result of direct optical
- 92 transitions with an energy of 1.42 eV.

Based on experimental results of investigation of temperature dependence of electrical conductivity and ranges of photocurrent of crystal  $InGaTe_2$  their power zonal chart represented in fig.1 is constructed.

- 96 In Fig.4 dependences of amplitude values of nonequilibrium photoconductivity of crystal  $InGaTe_2$ 97 on intensity of laser radiation are represented.
- As is seen from the figure 4, we observe considerable variety in the form of lux-Ampere charasteristics (LAC). In the case of excitement of low nuclear samples by the impulses of laser light with the wavelength of  $\lambda = 1,06$  mkm ( $\hbar \omega = 1,17 \text{eV}$ ), LAC at first has linear dependence, and at higher intensities the saturation has been observed (curve1). In high nuclear (pure) crystals, we observe square dependent (curve 2). The experiments carried out in high-nuclear samples of second harmonic Nd:YAG laser, ( $\lambda = 0,53$  mkm,  $\hbar \omega = 2,34 \text{eV}$ ) has the dependence I<sub>f</sub> ~I<sup>0,5</sup> (curve3).
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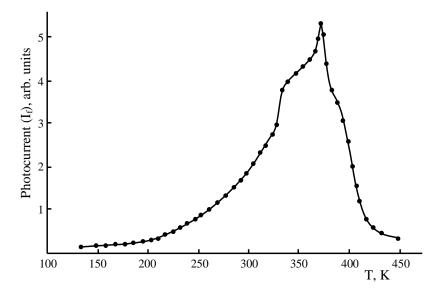


Fig. 3. Dependence of the photocurrent on the temperature in  $InGaTe_2$ 

The curves of relaxation of nonequilibrium photoconductivity for two groups of samples of excitation of the first (curve a,b) and second harmonic (curve c) Nd:YAG of the laser are given in Fig. 5. As is seen from the figure for low-nuclear samples, relaxation time is independent of light intensity, while for high-nuclear samples we observe its reduction with increase in power of a rating.

Before we pass to discussion of experimental results, it should be noted that by exciting semiconductors by ordinary light sources with the energy of quantum less than the width of the band structure ( $\hbar \omega < E_g$ ) we observe impurity photoconduction. In the case of impurity photoconductivity the Lux-amper characteristic (LAC) photoconductivity is of linear character and is saturated with growth of intensity of exciting light [11]. The last circumstance is stipulated by the fact that at rather great light intensities practically all carriers from impurity centers may be transfered by the light to the zone. Such a saturation of LAC is very seldom is observed at ordinarily used light intensities.

116 The obtained experimental results shows that in low-nuclear crystal  $InGaTe_2$  (E<sub>g</sub> =1,42eV) 117 excitement by laser radiation with energy of quantum  $\hbar\omega$  = 1,17eV the impurity photoconduction holds 118 (fig.4, curve 1).

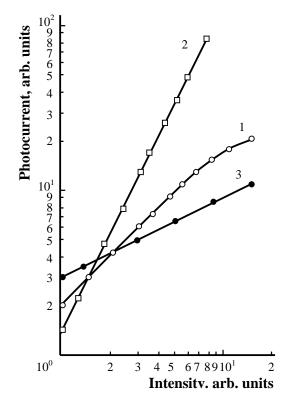


Fig.4. The dependence of the photocurrent on the laser intensity in  $InGaTe_2$ : 1, 2-low-resistance and high-resistance samples excited by the first harmonic of the laser, 3- high-resistance samples excited by the second harmonic of the laser.

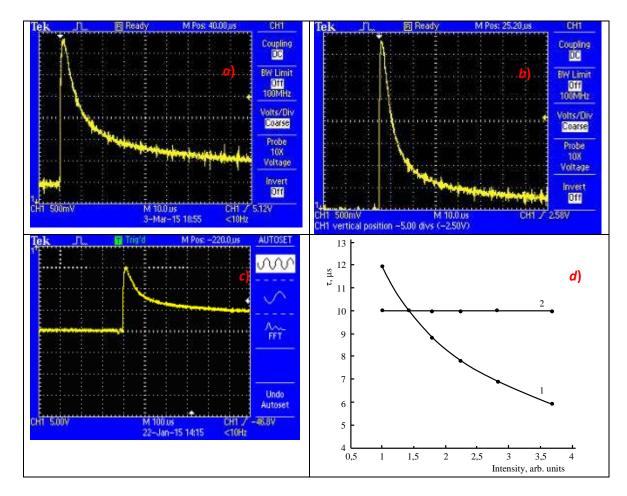
(2)

120 On the other hand, at excitement of the semiconductor by laser light with power of photon less 121 than the width of the band structure, the multiphoton light absorption in semiconductors corresponding to 122 birth or the couple electron-hole at simultaneous absorption of several photon [12,13] is also essential. There exist zone methods for investigating multiquantum transitions: optical absorption, investigation of 123 absorption of IK-light or extreme high-frequency radiation by nonequilibrium current carriers excited by 124 light source of high intensity, recombinational radiation photoconduction. All mentioned methods are 125 based directly or indirectly on the known relations between concentrations of nonequilibrium current 126 carriers  $\Delta n$  generated at absorption of n – photons in elementary act, by the light absorption coefficient  $\alpha$ , probability of multiquantum transition of W<sup>(n)</sup> and intensity of exciting light I: 127 128

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$$\alpha^n \sim I_0^{n-1}, \ \mathsf{W}^{(n)} \sim \Delta n \sim \mathsf{I}^{(n)}$$

0 In particular, at two-photon absorption 
$$\alpha^2 \sim I_{\odot}$$

130 In particular, at two-photon absorption  $\alpha^2 \sim I_0$ 



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Fig.5. The curve of relaxation of nonequilibrium photoconduction in  $InGaTe_2$  at laser excitation: low- nuclear (a) and high – nuclear (b) sample at excitement light with power of quantum –  $\hbar\omega = 1,17$  eV, c) high – nuclear sample at excitement light with power of  $\gamma$  quantum  $\hbar\omega = 2,34$  eV, d) dependence of relaxation time on intensity of excitement for high – nuclear (curve 1) and low – nuclear (curve 2) samples at excitement by laser radiation, with power of quantum  $\hbar\omega = 1,17$  eV.

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141 The most direct way of registering multiphoton absorption process is the study of nonequilibrium photoconductivity in semiconductors [14-16]. The considerable advantage of the method is that, here the 142 143 crystal width has no principal roll and therefore the most perfect crystals of small sizes may be used the 144 two - photon photoconduction may be detected, and consequently, the two - photon absorption 145 coefficient may be measured at intensity considerably less than the ones at which the two - point 146 absorption begins two occur in experiments on transmission therefore, the two - photon photoconduction, 147 may be a very useful method for detection of two - photon transition and a convenient way for measuring 148 the two-photon absorption coefficient supplementing the method of different measurement on 149 transmission.

- 150 As is known, photoconduction in semiconductors is determined by the formula
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 $\Delta \sigma = e \mu \alpha \beta I_0 \tau \tag{3}$ 

- 152 Where, e- is the elementary charge,  $\mu$  -is the mobility of carriers, a is the absorption coefficient,  $\beta$  is the 153 quantum fiend, I<sub>0</sub>-is light intensity, T- is the lifetime of nonequilibrium carriers.
- 154 At one photon absorption the absorption coefficient is independent of light intensity, so 155  $\Delta \sigma \sim I_0$ . At two – photon absorption is  $\alpha^2 \sim I_0$ , therefore  $\Delta \sigma \sim \alpha^2 \cdot I_0 \sim I_0^2$ . Existed of square

depended in Lav- Ampere characteristic of photoconduction testifies on two-photon photoconduction in crystal  $InGaTe_2$  under laser radiation. Comparison of the curves of relaxation of impurity and two – photon photoconduction (Fig.5 curves a, b) also shows this . At two- photon excitation, the relaxation time is much less (~2 times) in comparison with impurity photoconduction, and it depends on excitation level. By increasing the excitation intensity, the relaxation time decreases (fig. 5 d. curve 1) and here with the slow component of relaxation curve undergoes considerable charge. In the case of impurity photoconduction, the relaxation time is independent on excitation intensity (Fig. 5 d. curve 2).

163 It should be noted that at excitation of crystal  $InGaTe_2$  by the second harmonic of neodymium 164 laser, the power of the guantum ( $\hbar\omega = 2,34$  sB) considerably exceeds the width of the band structure of 165 the crystal (E<sub>g</sub>=1,42sB). Therefore, photoconduction is of bipolar character. In this case, we observe 166 square recombination that leads to the dependence  $\Delta \sigma(\Delta n) \sim \sqrt{I_0}$ . Experimental results confirm such

167 dependence  $\Delta \sigma$  to I<sub>0</sub> (fig.4, curve 3).

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### 4. CONCLUSION

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171 In monocrystal *p*- $InGaTe_2$  grown up by Bridgeman method at wide temperature range (200–500K), 172 electroconduction is stipulated by ionization deionization of acceptor centers forming energy level with 173 depth of occurance 0,2eV and 0,8eV in the band structure of crystal. The interzonal optical transitions in 174 monocrystals  $InGaTe_2$  is realized by indirect ( $E_g^{HII}$ =1,029B) and direct ( $E_g^{II}$ =1,429B) transitions. In the 175 range temperature 130-373K, with growth of temperature , the photosensitivity steadily increases. In 176 photogeneration of nonequilibrium current carriers by laser radiation, the mechanism of one-photon and 177 two-photon absorption of quantum is detected.

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