

Electrical and photoelectric properties of crystal $InGaTe_2$

ABSTRACT

The electroconductivity, photoconductivity and relaxation curves of the nonequilibrium photoconductivity in $InGaTe_2$ crystals under incandescent lamp light sources and laser light have been experimentally investigated. It is shown that the optical absorption in the $InGaTe_2$ due to indirect and direct optical transitions, the band gap of 1.02 eV and 1.42 eV, respectively. Acceptor centers are found, at a depth of 0.203 eV and 0.801 eV. At high levels optical excitation in $InGaTe_2$ observed two-photon absorption. It was found that the nonequilibrium photoconductivity observed in $InGaTe_2$ crystals under the action of 1st and 2nd harmonic of neodymium laser radiation, due to the impurity, bipolar and the two-photon photoconductivity

Keywords: compound $InGaTe_2$, indirect and direct optical transitions, non-linear optic properties, photoelectric properties

1. INTRODUCTION

Tellurides complex semiconductor compounds of the class $A^{III}B^{III}C_2^{VI}$ are of interest as materials for radiation detectors. Radiation resistance and effect of radiation on electrical conductivity is intensively studied in compounds as $TlGa(In)Te_2(Se,S)$ [1-6]. $InGaTe_2$ crystals in comparison with other compounds $A^{III}B^{III}C_2^{VI}$ little studied. The power zonal chart was calculated theoretically by the pseudopotential method [7] and dispersion curves of high frequency dielectric permittivity and temperature dependences of differential thermo-emf [8,9] were studied. The obtained data are not sufficient for making general idea on physical properties of $InGaTe_2$ monocrystals and as well as for specifying values of the fundamental parameters.

In this paper we report the results of the investigation of an experimental study of electroconductivity, photoconductivity in $InGaTe_2$ crystals at wide range of temperatures and at different levels of optical excitation.

As is known, at higher degrees of optic excitation there is high concentration nonequilibrium carrier. This caused achieved of new current certain qualitatively semiconductors with new properties - such as non-linear optical and photoelectric properties, effects of saturation in the absorption of light, the degeneracy of carriers. The use of laser radiation with tunable wavelength and with different excitation intensities can simultaneously detect impurity, its own as well as nonlinear photoconductivity, if any, are in the investigated crystals.

Experimental study of the generation-recombination processes in $InGaTe_2$ crystals by laser radiation is of particular interest due of the possibility of their application in optoelectronics.

2. METHODS OF CALCULATION

The single crystals $InGaTe_2$ were grown by the Bridgeman method have tetragonal crystal structure with spatial symmetry $D_{4h}^{18}(I4/mcm)$, with lattice parameters $a = 8,463 \text{ \AA}$ and $c = 6,981 \text{ \AA}$. The Bridgman method is a popular way of producing semiconductor crystals. The samples for measurements were cut out from an ingot in the form of a rectangle with the sizes of $6 \times 3 \times 1 \text{ mm}^3$ and lighting happens in the direction parallel to the crystallography axes "c". $InGaTe_2$ monocrystals had p-type conductivity and depending on the growing conditions had a resistivity in the range $(1 \cdot 10^3 \div 1 \cdot 10^6) \text{ Ohm} \cdot \text{cm}$. A silver paste is used as ohmic contact. **Photoconductivity $InGaTe_2$ crystals investigated incandescent lamp light sources in a stationary mode**, by the method of modulation of light intensity at frequency at 47kHz. At a high level of optical excitation, the pulse laser Nd: YAG with built-in generators of the second and third harmonics for generating radiation with a wavelength of 1064, 532, 355 nm, and tunable wavelength in the range of 410- 710 nm. The pulse width $\Delta t = 1 \cdot 10^{-8} \text{ sec}$ maximum power is $\sim 12 \text{ MVt/cm}^2$

47 Intensity of laser radiation was measured (changed) by means of calibrated neutral filters. In the work nonstationary
 48 digital system including memorable oscillograph and computer system was used. The experimental technigue is
 49 similar to one described in the paper [10].

50

51 3. RESULTS AND DISCUSSION

52

53 Dependence of electrical conductivity σ on temperature T is represented in Fig.1. As is seen from the figure, in the
 54 dependence $\lg \sigma \sim 10^3/T$ one can isolate two rectilinear areas belonging to generation of charge carrier with
 55 electronic transitions $E_1 = 0,203$ eV and $E_2 = 0,801$ eV. Samples have p-type conductivity, it should be noted
 56 that the temperature dependence of σ in the temperature range of 100-500 K, due to thermal ionization of acceptor
 57 centers and activation of filled electron acceptor. Temperature dependence of the electrical conductivity in different
 58 samples are identical.

59 Photoconductivity of $InGaTe_2$ crystals were investigated in the temperature range 132–373 K (Fig. 2).

60 Photocurrent range cover a very wide range of energy from 1,2 eV to 2,6 eV. Photosensitivity of the crystals
 61 steadily increases according to temperature increase. As is seen from the figure the long-wave edge of ranges of
 62 photocurrent don't subject to Moss criterion. The character optic transition in $InGaTe_2$ is seen from dependence

63 $I_f(h\nu)$. In the range $I_f^{1/2} \sim h\nu$ experimental points in the long-wave edge fall on a same line that indicates the
 64 presence of an indirect transfer of energy 1,02eV(Fig.2a).

65 By this method is quite justified determination width of the band structure of semiconductors. As absorption
 66 coefficient for indirect transitions is less than for direct transitions, its definition from the absorption range is not
 67 really a simple task. For that Many authors use more sensitive method as the photoconductivity method.

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It is known that the absorption coefficient for indirect transition for equals

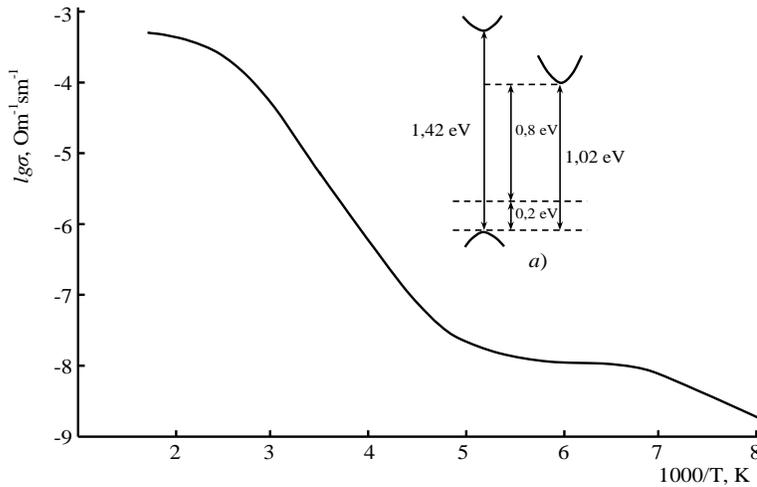


Fig. 1. The dependence of the electrical conductivity of crystal $InGaTe_2$ on the temperature. a) - the energy diagram $InGaTe_2$

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$$\alpha \sim (h\nu - E_g^{in} \pm E_p)^2 \quad (1),$$

70

71 where $h\nu$ - is the energy of quantum, E_g^{in} - is the width of the band structure, E_p -is the energy of the phonon that

72 participating in indirect optic transitions. As the photoconductivity size $I_f^{1/2} \sim \alpha^{1/2}$, then removing the dependence

73 of I_f on energy $h\nu$ in the weak abroption zone (at energies less than the direct zone) one can determine the width of

74

the band structure of indirect transitions.

75 Maximum photosensitivity of $InGaTe_2$ crystals is achieved at 1,42 eV. High photosensitivity in the
 76 fundamental absorption shows the presence of direct optical transitions in the crystal. Therefore, we can assume that
 77 in addition to indirect transition in the monocrystal $InGaTe_2$, a direct optical transition with energy 1,42eV causes
 78 the generation of nonequilibrium charge carriers.

79 Temperature dependence of photocurrent in $InGaTe_2$ is represented in Fig.3.
 80

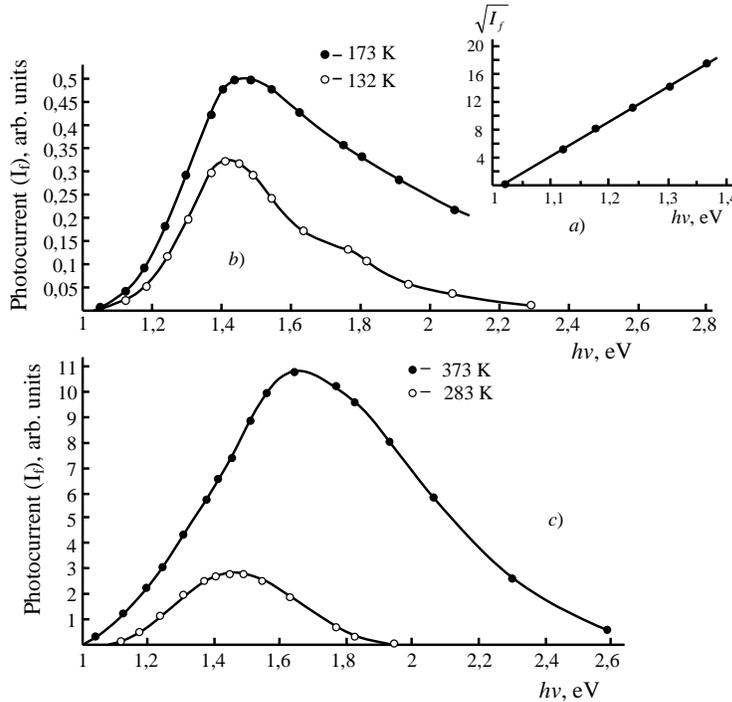


Fig. 2. Photocurrent spectra of $InGaTe_2$ crystals with different temperatures.

81 As seen from the Fig. 3, temperature quenching photocurrent happens at higher than 373 K. The range of
 82 photocurrent remained in the temperature clearing area has a sharp edge in the area of long-wave absorption and
 83 maximum of the range is at 1,42 eV. It is clear that in the area of temperature quenching of photocurrent main
 84 contribution to the photoconductivity make nonequilibrium charge carriers generated as a result indirect transitions
 85 with an energy of 1.02 eV and direct with an energy of 1.42 eV. Figure 1,a shows the energy band diagram
 86 $InGaTe_2$ crystals, built on the basis of experimental studies of the temperature dependence of the
 87 electroconductivity and photocurrent spectra.

88 In Fig.4 dependences of amplitude values of nonequilibrium photoconductivity of $InGaTe_2$ crystals on
 89 intensity of laser radiation are represented.

90 As seen from the figure 4, we observe considerable variety in the form of Lux-Ampere characteristics
 91 (LAC). In the case of excitement of low nuclear samples under impulses of laser light with the wavelength of $\lambda =$
 92 $1,06 \mu m$ ($\hbar\omega = 1,17 eV$) LAC at first has linear dependence and at higher intensities the saturation has been
 93 observed (curve1). In high nuclear (pure) crystals we observe square dependent (curve 2). The experiments carried
 94 out in high- nuclear samples of second harmonic Nd:YAG laser, ($\lambda = 0,53 \mu m$, $\hbar\omega = 2,34 eV$) has the dependence
 95 $I_f \sim I^{0,5}$ (curve3).
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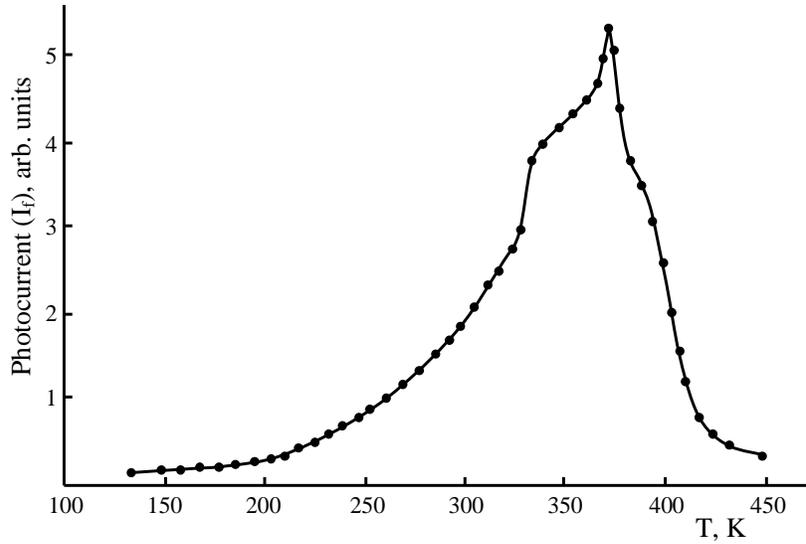


Fig. 3. Dependence of the photocurrent on the temperature in $InGaTe_2$

97 The curves of relaxation of nonequilibrium photoconductivity for two groups of samples of excitation of the first
 98 (curve a,b) and second harmonic (curve c) Nd:YAG of the laser are given in Fig. 5 .
 99 As seen from the figure, for the low-resistance samples relaxation time not depend on light intensity, while for high-
 100 nuclear samples we observe its reduction with increase in power of a rating.
 101 Before, we discuss the experimental results, it should be noted that the excitation of semiconductors incandescent
 102 lamp light sources, with a photon energy less than the band gap ($\hbar\omega < E_g$) observed impurity photoconductivity.
 103 In the case of impurity photoconductivity the Lux-Amper characteristic (LAC) photoconductivity is of linear
 104 character and is saturated with with increasing exciting light [11].
 105 The last circumstance is stipulated by the fact that at rather great light intensities practically all carriers from
 106 impurity centers may be transferred by the light to the zone. Such a saturation of LAC is very seldom is observed at
 107 ordinarily used light intensities. Our experimental results indicate that the low-resistivity $InGaTe_2$ crystals ($E_g =$
 108 1,42eV) when excited by laser radiation with a photon energy of $\hbar\omega = 1,17eV$ impurity photoconductivity appears
 109 (Figure 4, curve 1).
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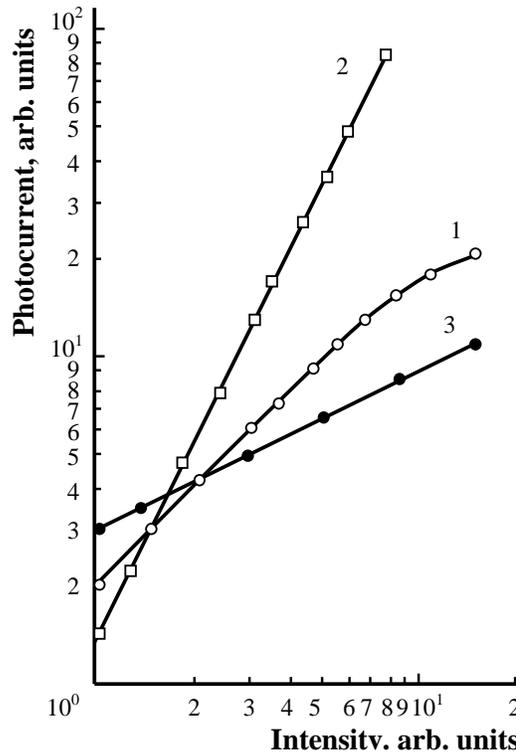


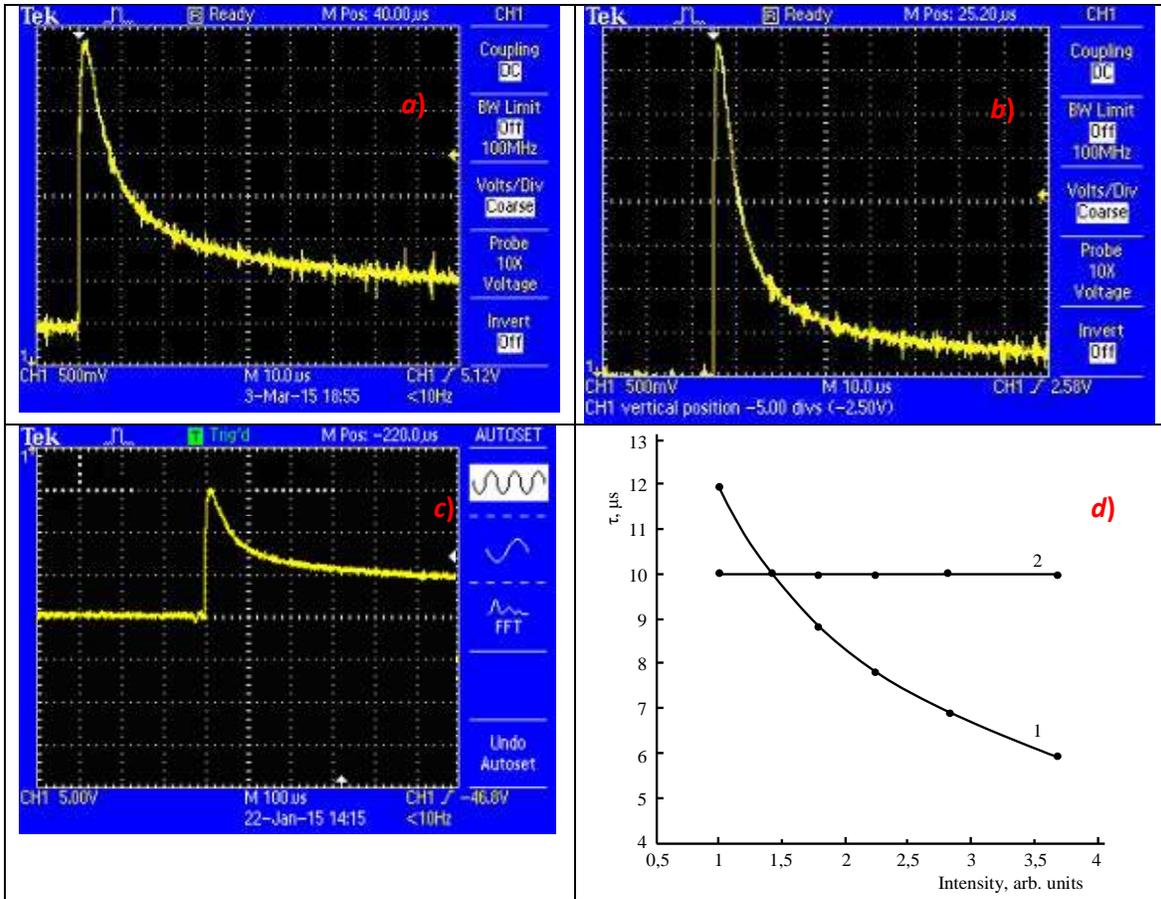
Fig.4. The dependence of the photocurrent on the laser intensity in $InGaTe_2$: 1, 2-low-resistance and high-resistance samples excited by the first harmonic of the laser, 3- high-resistance samples excited by the second harmonic of the laser.

112 At excitement of the semiconductor by laser light with power of photon less than the width of the band
 113 structure, the multiphoton light absorption in semiconductors corresponding to birth or the couple electron-hole at
 114 simultaneous absorption of several photon [12,13] is also essential. There exist zone methods for investigating
 115 multi-quantum transitions: optical absorption, investigation of absorption of IR-light or extreme high-frequency
 116 radiation by nonequilibrium current carriers excited by light source of high intensity, recombination radiation,
 117 photoconduction. All mentioned methods are based directly or indirectly on the known relations between
 118 concentrations of nonequilibrium current carriers Δn generated at absorption of n – photons in elementary act, by
 119 the light absorption coefficient α , probability of multi-quantum transition of $W^{(n)}$ and intensity of exciting light I :

$$120 \quad \alpha^n \sim I_0^{n-1}, \quad W^{(n)} \sim \Delta n \sim I^{(n)} \quad (2)$$

121 In particular, at two-photon absorption $\alpha^2 \sim I_0$.

122



123
 124
 125 **Fig.5. The curve of relaxation of nonequilibrium photoconduction in $InGaTe_2$ at**
 126 **laser excitation: low- nuclear (a) and high – nuclear (b) sample at excitement light with**
 127 **power of quantum – $\hbar\omega = 1,17$ eV, c) high – nuclear sample at excitement light with power**
 128 **of γ quantum $\hbar\omega = 2,34$ eV, d) dependence of relaxation time on intensity of excitement**
 129 **for high – nuclear (curve 1) and low – nuclear (curve 2) samples at excitement by laser**
 130 **radiation, with power of quantum $\hbar\omega = 1,17$ eV.**
 131
 132

133 The most direct way of registering multiphoton absorption process is the study of nonequilibrium photoconductivity
 134 in semiconductors [14-16]. The considerable advantage of the method is that, here the crystal width has no principal
 135 roll and therefore the most perfect crystals of small sizes may be used the two – photon photoconduction may be
 136 detected, and consequently, the two – photon absorption coefficient may be measured at intensity considerably less
 137 than the ones at which the two – point absorption begins two occur in experiments on transmission therefore, the two
 138 – photon photoconduction, may be a very useful method for detection of two – photon transition and a convenient
 139 way for measuring the two- photon absorption coefficient supplementing the method of different measurement on
 140 transmission.

141 As is known, photoconduction in semiconductors is determined by the formula

$$142 \quad \Delta\sigma = e\mu\alpha\beta I_0\tau \quad (3)$$

143 Where, e- is the elementary charge, μ - is the mobility of carriers, α is the absorption coefficient, β - is the quantum
 144 fiend, I_0 -is light intensity, τ - is the lifetime of nonequilibrium carriers.

145 At one – photon absorption the absorption coefficient is independent of light intensity, so $\Delta\sigma \sim I_0$. At two
 146 – photon absorption is $\alpha^2 \sim I_0$, therefore $\Delta\sigma \sim \alpha^2 \cdot I_0 \sim I_0^2$. Existed of square depended in Lux-Ampere

147 characteristic of photoconduction testifies on two-photon photoconduction in $InGaTe_2$ crystals under laser radiation.
148 Comparison of the curves of relaxation of impurity and two – photon photoconduction (Fig.5 curves a, b) also shows
149 this. At two- photon excitation, the relaxation time is much less (\sim in 2 times) in comparison with impurity
150 photoconduction, and it depends on excitation level. By increasing the excitation intensity, the relaxation time
151 decreases (Fig. 5 d. curve 1) and here with the slow component of relaxation curve undergoes considerable change.
152 In the case of impurity photoconduction, the relaxation time is independent on excitation intensity (Fig. 5 d. curve 2).

153 It should be noted that at excitation of $InGaTe_2$ crystals by the second harmonic of neodymium laser, the
154 photon energy ($\hbar\omega = 2,34$ eV) is much higher than the band gap of the crystal ($E_g=1,42$ eV). Therefore,
155 photoconduction is of bipolar character. In this case, we observe square recombination that leads to the dependence
156 $\Delta\sigma(\Delta n) \sim \sqrt{I_0}$. Experimental results confirm such dependence $\Delta\sigma$ to I_0 (Fig.4, curve 3).

157

158 4. CONCLUSION

159

160 Monocrystal $p-InGaTe_2$ grown up by Bridgeman method at wide temperature range (200–500K),
161 electroconduction is stipulated by ionization deionization of acceptor centers forming energy level with depth of
162 occurrence 0,2eV and 0,8eV in the band structure of crystal. The interzonal optical transitions in $InGaTe_2$
163 monocrystals is realized by indirect ($E_g^{ind}=1,02\text{eV}$) and direct ($E_g^{dir}=1,42\text{eV}$) transitions. In the range temperature
164 130-373K, with growth of temperature, the photosensitivity steadily increases. In photogeneration of nonequilibrium
165 current carriers by laser radiation, the mechanism of one-photon and two-photon absorption of quantum is detected.
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